

Функциональный интеграл
для многочастичной КМ

$$\hat{H} = \int \left(\frac{\nabla \hat{\psi}^\dagger \nabla \hat{\psi}}{2m} + U(x) \hat{\psi}^\dagger(x) \hat{\psi}(x) \right) dx + \frac{1}{2} \int d\vec{x} d\vec{x}' \hat{\psi}^\dagger(x) \hat{\psi}^\dagger(x') U(x-x') \hat{\psi}(x') \hat{\psi}(x)$$

$$\Leftrightarrow H_N = \sum_i \left(\frac{p_i^2}{2m} + U(x_i) \right) + \frac{1}{2} \sum_{i \neq j} U(x_i - x_j)$$

Базис когерентных состояний:

$$1) [\hat{a}_i, \hat{a}_j^\dagger] = \delta_{ij} \quad \text{Бозоны}$$

$$\hat{a} |\varphi\rangle = \varphi |\varphi\rangle$$

$$|\varphi\rangle = e^{\varphi \hat{a}^\dagger} |0\rangle = \sum_n \frac{\varphi^n}{n!} |n\rangle$$

$$\langle 0 | \varphi \rangle = e^{-\varphi^2/2}$$

$$\hat{I} = \sum_n |n\rangle \langle n| = \int \frac{d(\operatorname{Re} \varphi) d(\operatorname{Im} \varphi)}{\pi} e^{-|\varphi|^2} |\varphi\rangle \langle \varphi|$$

2) Фермионы

$$\{\hat{a}_i, \hat{a}_j^\dagger\} = \delta_{ij}$$

$$\hat{a}_i |q_i\rangle = q_i |q_i\rangle \quad \hat{a}_i^2 |q_i\rangle = 0$$

$$\hat{a}_j |q_j\rangle = q_j |q_j\rangle \quad \hat{a}_j^2 |q_j\rangle = 0$$

$$\{q_i, q_j + q_j q_i\} = (q_i q_j + q_j q_i) |q_i, q_j\rangle$$

Грассмановы числа

1) $q_i q_j + q_j q_i = 0$
 $\Rightarrow q^2 = 0$

2) $(q_i q_j) q_k = q_i (q_j q_k)$
 $q_i (q_j + q_j q_k) = q_i q_j + q_i q_j q_k$

3) $\{q_i\}_{i=1}^N \Leftrightarrow$ вектор $q = \sum_{i=1}^N a_i q_i + \sum_{i,j=1}^N a_{ij} q_i q_j + \dots + a_{12\dots N} q_1 \dots q_N$

4) $f(x+q) = f(x) + f(q)$

5) $\frac{\partial}{\partial q_i} q_j = \delta_{ij}$ $\frac{\partial}{\partial q} (S(q)) = \frac{\partial}{\partial q} (q f)$

6) $\int d q_i q_i = \delta_{ii}$ $\int d q_i q_i q_i = -\int d q_i q_i = -1$

7) "Комплексное сопряжение"

$$q \rightarrow \bar{q} \quad \bar{\bar{q}} = q$$

8) Грассмановы интегралы

$$\int \prod_i (d\bar{q}_i dq_i) \exp(-\sum_{ij} \bar{q}_i L_{ij} q_j) = \det L$$

Аналог: $\int \prod_i d\bar{q}_i dq_i \exp(-\sum_i \lambda_i \bar{q}_i q_i)$
 $\prod_i (1 - \lambda_i \bar{q}_i q_i) = \int \prod_i [d\bar{q}_i dq_i (\delta - \lambda_i \bar{q}_i q_i)] = \prod_i \lambda_i = \det L$

Грассмановы числа, аналог

Fermi: $\hat{a} |q\rangle = q |q\rangle$
 $|q\rangle = e^{\hat{a}^\dagger q} |0\rangle = (1 + \hat{a}^\dagger q) |0\rangle = |0\rangle + q \hat{a}^\dagger |0\rangle$
 $\langle q | s \rangle = (\langle 0 | + \bar{q} \langle 1 |) (|0\rangle + q |1\rangle) = 1 + \bar{q} q = e^{\bar{q} q}$
 $\hat{I} = |0\rangle \langle 0| + |1\rangle \langle 1| = \int d\bar{q} dq e^{-\bar{q} q} |q\rangle \langle q|$
 $= \int d\bar{q} dq (1 - \bar{q} q) (|0\rangle + q |1\rangle) (\langle 0| + \bar{q} \langle 1|) = |0\rangle \langle 0| + |1\rangle \langle 1|$

Path integral для КМ

1) $\hat{H} = \omega \hat{a}^\dagger \hat{a}$

$$\langle \varphi' | (1 - i\hat{H}\epsilon) | \varphi \rangle = e^{\bar{\varphi}' \varphi - iH(\bar{\varphi}', \varphi) \epsilon} = e^{\bar{\varphi}' \varphi} e^{-iH(\bar{\varphi}', \varphi) \epsilon}$$

$$\langle \varphi' | \hat{H}(\hat{a}^\dagger, \hat{a}) | \varphi \rangle = H(\hat{a}^\dagger \rightarrow \bar{\varphi}', \hat{a} \rightarrow \varphi) \langle \varphi' | \varphi \rangle$$

$$\hat{I} = \int \prod_{k=1}^N d\bar{\varphi}_k d\varphi_k \exp(-\sum_{k=1}^N \bar{\varphi}_k \varphi_k + \sum_{k=1}^N \bar{\varphi}_k \varphi_{k+1} - i\epsilon \sum_{k=1}^N H(\bar{\varphi}_k, \varphi_{k+1})) = \int \mathcal{D}[\bar{\varphi}(t), \varphi(t)] \exp(iS[\bar{\varphi}(t), \varphi(t)])$$

$$S = \int dt (i\bar{\varphi} \partial_t \varphi - H(\bar{\varphi}, \varphi))$$

2) $\hat{H} = \sum_i \hat{a}_i^\dagger \hat{a}_i \cdot E_i \Leftrightarrow \int \left(\frac{\nabla \psi^\dagger \nabla \psi}{2m} + U(x) \psi^\dagger \psi \right) dx$
 $\hat{\psi}(x) = \sum_i \varphi_i^M \cdot \hat{a}_i$

$$\Rightarrow \langle \Omega | T \{ \hat{\psi}(\vec{x}, t) \hat{\psi}^\dagger(\vec{y}, t') \} | \Omega \rangle = \frac{\int \mathcal{D}[\bar{\psi}, \psi] \cdot \psi(\vec{x}, t) \bar{\psi}(\vec{y}, t') e^{iS[\bar{\psi}, \psi]}}{\int \mathcal{D}[\bar{\psi}, \psi] e^{iS[\bar{\psi}, \psi]}} \quad \text{Fermi, Bose}$$

$$H[\bar{\psi}, \psi] = \int d\vec{x} d\vec{x}' \left(\frac{\nabla \bar{\psi} \nabla \psi}{2m} + U(\vec{x} - \vec{x}') \bar{\psi}(\vec{x}) \bar{\psi}(\vec{x}') \right)$$

$$S = \left[\int dt dx (i\bar{\psi} \partial_t \psi) - H \right]$$

Самобольеупадающий бозе-газ

N-частичный бозе-газ

$$\langle \Psi | \int d^3x \hat{\psi}^\dagger(x) \hat{\psi}(x) | \Psi \rangle = N$$

$$\Leftrightarrow \text{функция нуклеона} \quad \int d^3x |\Psi|^2 = N$$

Остаток $= \min_{\langle \Psi | \hat{H} | \Psi \rangle = N} \langle \Psi | \hat{H} | \Psi \rangle = \min_{\langle \Psi | \hat{H} | \Psi \rangle = N} (\langle \hat{H} \rangle - \mu \langle N \rangle) \rightarrow |\Psi(\mu)\rangle$
 или потенциал μ найдет из уравнения $\langle \Psi(\mu) | \hat{N} | \Psi(\mu) \rangle = N$

$\hat{H} \mapsto \hat{H}' = \hat{H} - \mu \hat{N}$, но без граничных условий.

Большой канонический ансамбль

N бозе-газ

$$U(x-y) = g \delta(x-y) \quad g - \text{свободная энергия}$$

$$\hat{H}' = \int \left(\frac{\nabla \psi^\dagger \nabla \psi}{2m} + \mu \psi^\dagger \psi + \frac{g}{2} \psi^\dagger \psi^\dagger \psi \psi \right) dx$$

$$Z = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp \left(i \int dt dx \left(i\bar{\psi} \partial_t \psi - \frac{|\nabla \psi|^2}{2m} + \mu \psi^\dagger \psi - \frac{g}{2} |\psi|^4 \right) \right)$$

"метод перевала"

$$\frac{\delta S}{\delta \psi} = 0 \Rightarrow \left(i \partial_t \psi + \frac{\nabla^2 \psi}{2m} + \mu \psi - g |\psi|^2 \psi = 0 \right)$$

$\psi(\vec{x}, t) = \text{const}$ $\mu \psi = g |\psi|^2 \psi$ $g > 0$ - отталкивание

$\int |\psi|^2 dx = \text{const}$ $|\psi|^2 = \mu/g \Rightarrow \psi_0 = \sqrt{\mu/g} e^{i\varphi}$ $\delta \text{бозе}$ $- \text{конденсат}$

$\psi = \psi_0 + \delta\psi$ спонтанное нарушение симметрии $U(1)$

$$0 = \langle \hat{\psi}(\vec{x}, t) \rangle = \frac{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi \psi e^{iS}}{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS}} = \psi_0$$

$\hat{N} |\Omega\rangle = N |\Omega\rangle$

$$\langle \Omega | \hat{\psi}(\vec{x}, t) \hat{\psi}^\dagger(\vec{x}', t') | \Omega \rangle = |\psi_0|^2 + \text{погр.}$$

Дальний порядок $|\vec{x} - \vec{x}'| \rightarrow \infty$

$|\Omega\rangle$ - когерентное состояние

приближение $\Rightarrow \hat{\psi} |\Omega\rangle = \psi_0 |\Omega\rangle$

$\hat{\psi} |\Omega_1\rangle = \sqrt{\frac{N}{g}} e^{i\varphi_1} |\Omega_1\rangle$
 $\hat{\psi} |\Omega_2\rangle = \sqrt{\frac{N}{g}} e^{i\varphi_2} |\Omega_2\rangle$

$$\langle \Omega_1 | \Omega_2 \rangle = \exp \left(- \int dx \bar{\psi}_1 \psi_2 \right) = \exp \left(- \frac{\mu}{g} \cdot V e^{i(\varphi_1 - \varphi_2)} \right)$$

$$P = \frac{|\langle \Omega_1 | \Omega_2 \rangle|^2}{|\langle \Omega_1 | \Omega_1 \rangle| \cdot |\langle \Omega_2 | \Omega_2 \rangle|} = \exp \left(- 2 \frac{\mu}{g} V (1 - \cos \Delta\varphi) \right)$$

10^{23} $\Delta\varphi = O(1)$