

Формализм Гельфанда-Явновского

$$\frac{\det(\hat{H}_S)}{\det(\hat{H}_0)} = \prod_n \left(\frac{\lambda_n^{(S)}}{\lambda_n^{(0)}} \right) \quad H_S = -\partial_x^2 + \frac{s(s+1)}{ch^2 x} \quad H_0 = -\partial_x^2$$

I способ (Нумерное представление)

$\epsilon = L/N$

$\hat{H} = -\partial_x^2 + U(x)$

$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$

$H \rightarrow \text{mat}_{(N-1) \times (N-1)}$

$(\hat{H}\psi)_k = -\frac{\psi_{k-1} - 2\psi_k + \psi_{k+1}}{\epsilon^2} + U(x_k)\psi_k$

$\hat{H} = \begin{pmatrix} U(x_1) + \frac{2}{\epsilon^2} & -\frac{1}{\epsilon^2} & & \\ -\frac{1}{\epsilon^2} & U(x_2) + \frac{2}{\epsilon^2} & -\frac{1}{\epsilon^2} & \\ & -\frac{1}{\epsilon^2} & \ddots & \\ & & & U(x_{N-1}) + \frac{2}{\epsilon^2} \end{pmatrix}$

$\hat{H} = \frac{1}{\epsilon^2} \begin{pmatrix} \epsilon^2 U(x_1) + 2 & -1 & & \\ & -1 & \epsilon^2 U(x_2) + 2 & -1 \\ & & \ddots & \ddots \\ & & & -1 & \epsilon^2 U(x_{N-1}) + 2 \end{pmatrix}$

$f_{N-1} = \det(\dots)$

$\psi_0 = 0$

$\psi_N = 0$

$\det \hat{H} = \frac{1}{\epsilon^{N-1}} f_{N-1}$

$f_{N-1} = (U(x_{N-1})\epsilon^2 + 2)f_{N-2} - (-1) \times (-1)f_{N-3}$

$f_0 = 0$

$f_1 = 2 + \epsilon^2 U(x_1)$

$f'(0) = \frac{f_1 - f_0}{\epsilon} = \frac{2}{\epsilon}$

$\epsilon \rightarrow 0$; $f_\epsilon \equiv f(\epsilon) \Rightarrow -f'' + U(x)f(x) = 0$

$1) \det \equiv f(L) = \frac{2}{\epsilon} g(L)$

$2) \begin{cases} g(0) = 0 \\ g'(0) = 1 \end{cases} \Rightarrow \det \hat{H} = \frac{1}{\epsilon^{N-1}} g(L) \Rightarrow \frac{\det(H_1)}{\det(H_2)} = \frac{\det(-\partial_x^2 + U_1(x))}{\det(-\partial_x^2 + U_2(x))} = \frac{g_1(L)}{g_2(L)}$

II способ

$\frac{\det(H_1 - E)}{\det(H_2 - E)} = \frac{f_1(E)}{f_2(E)}$

$f_i(E) = \prod_n (\lambda_n^{(i)} - E)$

$f_i(E) = 0 \Leftrightarrow E$ - собственное значение

Прогнать функцию, к-я равна 0 $\Leftrightarrow$ собственное значение

Метод стрелок

$f(0) = 0$

$f'(0) = 1$

$f(L|E)$

$f(L|E_n) = 0$

$f(x|E_n)$ - собственная функция

$f(L|E) \equiv g(E)$ - характеристический полином

(L тогда же го членом ф-ции)

Если $f_i(x \rightarrow 0)$ сходится асимпт. $\Rightarrow$ член ф-ции сгуст

Примеры

1) Квантовый осциллятор

$S[x] = \int_0^\beta d\tau ((\partial_\tau x)^2 + \omega^2 x^2)$

$\hat{H}_\omega = -\partial_\tau^2 + \omega^2$

$\frac{\det(\hat{H}_\omega)}{\det(\hat{H}_0)} = \frac{\text{sh}(\omega\beta)/\omega}{\beta} = \frac{\text{sh}(\omega\beta)}{\omega\beta}$

$(-\partial_\tau^2 + \omega^2)f(\tau) = 0$

$f(\tau) = C_1 \text{ch}(\omega\tau) + C_2 \text{sh}(\omega\tau)$

$-\partial_\tau^2 f_2(\tau) = 0$

$f_2(\tau) = C_1 + C_2 \tau$

$f(0) = 0 \Rightarrow C_1 = 0$

$f(x=0) = \tau \Rightarrow C_2 = \frac{1}{\omega}$

$\Rightarrow f(\tau) = \frac{\text{sh}(\omega\tau)}{\omega}$

2) Унстаптон "E = -2e"

$\frac{\det(H_S + 2e^2)}{\det(H_0 + 2e^2)}$

$H_S = -\partial_\tau^2 + \frac{s(s+1)}{ch^2 \tau}$

Для ф-ции $s=2$

1) $L \rightarrow \infty \Rightarrow$ промат в левую $2L \rightarrow \infty$

2) Потенциал $z\text{e}^{2x}$

$\hat{H} = \begin{pmatrix} z\text{e}^{2x} & 0 \\ 0 & H_0 + 2z\text{e}^{2x} \end{pmatrix}$

$\Rightarrow \det \hat{H} = \det(\hat{H}^{(+)}) \det(\hat{H}^{(-)})$

$\det(\hat{H}^{(+)}) = \psi^{(+)}(L)$

$\psi^{(+)}(L) \rightarrow z\text{e}^{2x}$

$\begin{cases} \psi_0(0) = 1 \\ \psi_0'(0) = 0 \end{cases}$

$\det(\hat{H}^{(-)}) = \psi^{(-)}(L)$

$\begin{cases} \psi^{(-)}(0) = 0 \\ \psi^{(-)}'(0) = 1 \end{cases} \quad x \leftrightarrow \tau$

$-\psi'' + \frac{s(s+1)\psi(x)}{ch^2 x} = -2e^2 \psi(x)$

$y = ch^2 x \Rightarrow y(1-y)\psi'(y) + (\frac{1}{2}-y)\psi(y) - (\frac{s(s+1)}{4y} - \frac{e^2}{4})\psi(y) = 0$

$\psi(y) = y^{s+1/2} f(y) \Rightarrow y(1-y)y''(y) + [(s+3/2)-y(s+1)]f'(y) - \frac{(s+1)^2 - 2e^2}{4}f(y) = 0$

$z = 1-y = -sh^2 x \Rightarrow z(1-z)y''(z) + (\frac{1}{2}-z(s+1))f'(z) - \frac{(s+1)^2 - 2e^2}{4}f(z) = 0$

$\begin{cases} a = \frac{s+1+2e}{2} \\ b = \frac{s+1-2e}{2} \\ c = 1/2 \end{cases} \Rightarrow {}_2F_1(a, b; c; z) \sim {}_2F_1(b-c+1, a-c+1; 2-c; z)$

$\begin{cases} \psi^{(+)}(x) = \text{ch}^{s+1/2} x \cdot {}_2F_1(\frac{s+1+2e}{2}, \frac{s+1-2e}{2}; \frac{1}{2}; -sh^2 x) \\ \psi^{(-)}(x) = sh^x \cdot \text{ch}^{s+1/2} x \cdot {}_2F_1(\frac{s-2e+2}{2}, \frac{s+2e+2}{2}; \frac{3}{2}; -sh^2 x) \end{cases} \quad E = -2e^2$

$\Rightarrow \frac{\det(H_S + 2e^2)}{\det(H_0 + 2e^2)} = \frac{\psi_0^{(+)}(L)\psi_0^{(-)}(L)}{\psi_0^{(+)}(L)\psi_0^{(-)}(L)} \quad (L \rightarrow \infty)$

$\begin{cases} \psi_0^{(+)}(x) = \frac{\Gamma(1/2)\Gamma(2e)}{2^{2e}\Gamma(\frac{1+s+2e}{2})\Gamma(\frac{2e-s}{2})} e^{2ex} \\ \psi_0^{(-)}(x) = \frac{\Gamma(1/2)\Gamma(2e)}{2^{2e}\Gamma(\frac{s+2e+2}{2})\Gamma(\frac{1-s+2e}{2})} e^{2ex} \end{cases} \quad x \rightarrow \infty$

$\begin{cases} \psi_0^{(+)} = \text{ch}^{2e} x \sim \frac{1}{2} e^{2ex} \\ \psi_0^{(-)} = \frac{sh(2ex)}{2e} \sim \frac{1}{2e} e^{2ex} \end{cases}$

$\frac{\det(H_S + 2e^2)}{\det(H_0 + 2e^2)} = \frac{2e \Gamma^2(2e)}{\Gamma(1+s+2e)\Gamma(2e-s)}$

Примеры

$\det'(-\partial_\tau^2 + \omega^2 - \frac{3\omega^2}{2ch^2 \tau}) = \frac{\det(-\partial_\tau^2 + 4 - \frac{6}{ch^2 x})}{\det(-\partial_\tau^2 + \omega^2)}$

$x=2; s=2$

Нумерное мога!

$x^2 = 4 + \epsilon \Rightarrow$ нумерное мога $\rightarrow \epsilon$

$\Gamma(x) \sim \frac{1}{x}$

$\det'(H_S + 2e^2) = \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} \det(H_S + 2e^2 + \epsilon)$

$2e = 2 + \frac{\epsilon}{4} \Rightarrow \frac{\det'(\dots)}{\det'(\dots)} = \lim_{\epsilon \rightarrow 0} \frac{\frac{1}{\epsilon} \det'(\dots)}{\frac{1}{4} \det'(\dots)} = \frac{2 \Gamma^2(2)}{\Gamma(5) \cdot \Gamma(\frac{5}{4})} \cdot \frac{4}{\epsilon} = \frac{2}{24} = \boxed{\frac{1}{12}}$